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WIMPs with mass ~ 100 GeV from a pair of massless and small-mass scalar-fields in interaction

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Abstract. – It is shown that a massless scalar-field can interact with a scalar-field bearing particles of mass m to produce localized particlelike concentrations of field energy, WIMP *solitons* with a mass M orders of magnitude greater than m . The characterizing signature features of such solitons are described, and their possible detection by ongoing observations is noted.

Tentative experimental evidence for WIMPs, Weak Interacting Massive Particles with a mass of the order 100 GeV, has been reported recently [1,2]. Conjectured to be the principal component of cold dark matter and hence about 84% of all matter, WIMPs are in striking contrast to the near-massless particles of previously considered scalar-field candidates for dark matter, fields which carry masses that range from $\sim 10^{-22}$ eV [3] to $\sim 10^{-3}$ eV [4]. Remarkably however, a pair of massless and small-mass scalar-fields in interaction can produce WIMPlike entities with a mass of the order 100 GeV, as shown below. Since observational evidence for self-interacting cold dark matter has been established for over a decade [5], this model for WIMPs is a viable alternative to a supersymmetry particle with WIMP properties.

Now after a century of particle-theory development, there has emerged a very large order-of-magnitude range in seemingly fundamental masses [6,7]. We have the hypothetical Planck mass $G^{-1/2} = 1.22 \times 10^{19}$ GeV [8], the (probably) massless photon ($m_\gamma \lesssim 2 \times 10^{-11}$ eV), the lightest mass-state neutrino ($m_{\nu_1} \lesssim 5 \times 10^{-3}$ eV), and finally the suspect-candidate dark matter particles, conjectured to be as large as ~ 100 GeV [1,2], as small as $\sim 10^{-3}$ eV [4], or perhaps even effectively massless at $\sim 10^{-22}$ eV [3]. Nestled in between and spanning five orders-of-magnitude from the weak-interaction gauge bosons with masses ~ 100 GeV down to the electron with

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$m_e = 0.511 \text{ MeV}$, we have the experimentally established particles. In a future fundamental theory, the masses which enter must surely bear basic theoretical relationships to one another. Hence, considerable interest is attached to field-theoretic mechanisms that may account for order-of-magnitude differences in observable particle mass. The purpose of this communication is to describe such a mechanism, which may in fact apply to WIMPs.

Consider the effective (post-renormalization) Lagrangian

$$\mathcal{L}_{eff} = -\frac{1}{2} g^{\mu\nu} (\chi_{,\mu} \chi_{,\nu} + \phi_{,\mu} \phi_{,\nu}) - \frac{1}{2} m^2 \chi^2 + \beta \chi \phi^3 \quad (1)$$

in which χ is a real scalar-field bearing particles of mass m , ϕ is a massless real scalar-field, and β is a dimensionless coupling-constant. From (1) we obtain the field equations

$$\begin{aligned} g^{\mu\nu} \chi_{,\mu\nu} - m^2 \chi + \beta \phi^3 &= 0 \\ g^{\mu\nu} \phi_{,\mu\nu} + 3\beta \chi \phi^2 &= 0 \end{aligned} \quad (2)$$

and the canonical energy density

$$T_{00} = \frac{1}{2} \left((\chi_{,0})^2 + |\nabla \chi|^2 + (\phi_{,0})^2 + |\nabla \phi|^2 \right) + \frac{1}{2} m^2 \chi^2 - \beta \phi^3 \chi \quad (3)$$

Let us suppose that Eqs. (2) admit a static solution with $\chi_{,0} = 0 = \phi_{,0}$ in an appropriate Lorentz frame, and let us seek a singularity-free solution with spherical symmetry about the point $x \equiv (x_1, x_2, x_3) = 0$. The first member of (2) can be recast as the integral equation

$$\chi(|x|) = \frac{\beta}{4\pi} \int \frac{(\exp - m|x-y|)}{|x-y|} \phi^3(|y|) d^3 y \cong \beta m^{-2} \phi^3(|x|) \quad (4)$$

where the last member of (4) follows approximately if and only if

$$|\phi(|x|)/\phi(0)|^3 \ll 1 \quad \text{for} \quad |x| \gtrsim m^{-1} \quad (5)$$

i.e., if and only if $|\phi(|x|)/\phi(0)|$ declines from 1 to near 0 as a $|x|$ increases from 0 to become greater than m^{-1} .

Then, with substitution of the last member of (4) into the second member of (2), we get the elliptic partial differential equation in the ϕ field exclusively,

$$\nabla^2 \phi + 3\beta^2 m^{-2} \phi^5 = 0 \quad (6)$$

Remarkably, (6) admits the exact rigorous spherically-symmetric singularity-free *soliton* solution

$$\phi(|x|) = \pm \left(a|x|^2 + a^{-1} \beta^2 m^{-2} \right)^{-1/2} \quad (7)$$

in which a is a positive (dimensionless and disposable) constant of integration. We now verify *a posteriori* that for suitable values of a the solution (7) satisfies the approximation requirement (5):

$$\left| \phi(|x|)/\phi(0) \right|^3 = \left(a^2 \beta^{-2} m^2 |x|^2 + 1 \right)^{-3/2} \ll 1 \quad \text{for} \quad |x| \gtrsim m^{-1} \quad (8)$$

Indeed, (8) holds if and only if

$$|\beta| \ll a \quad . \quad (9)$$

Since the approximate radial size of the solution (7) is given by $\hat{r} = a^{-1} |\beta| m^{-1}$, (9) states that radial size is sub-Compton, *i.e.*, small compared to m^{-1} . The total field energy mass of the soliton (7) follows from (3) and (4) as

$$M \equiv \int T_{00} d^3x = \pi^2 m/4 |\beta| \quad (10)$$

with terms combining inside the integral, the subsequent integration being exact, and the constant of integration a scaling out. Hence, for $|\beta| \ll 1$ we have obtained solitons with a mass $M \gg m$ (while conversely, for a superstrong coupling with $|\beta| \gg 1$, we obtain $M \ll m$). It can be shown by long-established classical-field analysis that for all magnitudes of $|\beta|$ and a these solitons are anti-Coulombic, with like-sign solutions (7) attracting each other and unlike-sign solutions repelling each other [9]. The mass relation (10), the variable radial size feature of the solitons (with a disposable and $\hat{r} = \pi^2/4aM$) and their anti-Coulombic character may provide an identifying signature for them. In particular, if dark matter WIMPs with $M \sim 100$ GeV interact with terrestrial matter to afford detection by the most advanced signature-determining *CDMS II* apparatus [10], and if they show the latter soliton signature features, then such WIMPs may derive from interaction between a pair of massless and small-mass dark matter scalar-fields, such as those conjectured on the basis of cosmological observations [3,4].

In summary, it has been shown that a massless scalar-field can interact with a scalar-field bearing particles of mass m to produce localized particlelike concentrations of field energy, *solitons* with a mass M orders of magnitude greater than m . The characterizing signature features of such solitons have been described, and such signature features may show up in the ongoing *CDMS II* observations [10]. If so, WIMPs with $M \sim 100$ GeV may be solitons that follow from the χ and ϕ fields, as derived rigorously here from the effective Lagrangian (1).

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